# SUPPLEMENTARY INFORMATION 

# Structural characterization of ice XIX as the second polymorph related to ice VI 

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## Supplementary Discussion

Let's start with the calorimetric analysis of the second endotherm (ice XV $\rightarrow$ ice VI). Apparently, the second endotherm is much larger ( $48 \mathrm{~J} \mathrm{~mol}^{-1}$ ) for the top two traces pertaining to $\mathrm{H}_{2} \mathrm{O}$ ice $\beta$-XV, compared to all traces between 0.5 and $30 \% \mathrm{H}_{2} \mathrm{O}\left(7 \mathrm{~J} \mathrm{~mol}^{-1}\right)$, in which $\mathrm{D}_{2} \mathrm{O}$ ice $\beta-\mathrm{XV}$ forms at 1.8 GPa. This prompts the question why the ice XV to ice VI transition is suppressed in $\mathrm{D}_{2} \mathrm{O}$ ice $\beta$-XV samples. The answer is kinetics. In order to transform one ordered phase to another differently ordered phase an activation barrier needs to be overcome, where the transition state is necessarily disordered. We call this transition state ice $\mathrm{VI}^{\ddagger}$. This transition state then needs to order again, and turn into ice XV . Ordering of ice $\mathrm{VI}^{\ddagger}$ is much faster for $\mathrm{H}_{2} \mathrm{O}$ samples than for $\mathrm{D}_{2} \mathrm{O}$ samples. The large kinetic isotope effect is evident from the width of the first endotherm, which increases from $12 \pm 1 \mathrm{~K}$ in protiated samples to $18 \pm 2 \mathrm{~K}$ in deuterated ones (see blue double-arrows in Figure 1c). In other words, the 2.5 minutes provided for the order in ice XV to develop inside the calorimetry instrument at a scanning rate of $10 \mathrm{~K} / \mathrm{min}$ are not sufficient in deuterated samples, but in protiated ones. In order to make this point more evident we have conducted a calorimetry experiment, in which the heating scan was paused for 30 minutes at 120 K . In these 30 minutes there is ample time for $\mathrm{D}_{2} \mathrm{O}$ ice XV to form, and upon commencing the heating scan a large second endotherm of $70 \mathrm{~J} \mathrm{~mol}^{-1}$ is also observed for $\mathrm{D}_{2} \mathrm{O}$ samples (see dashed green line in Figure 1a at $5 \% \mathrm{H}_{2} \mathrm{O}$ ). That is, 30 minutes at 120 K are enough for deuterated ice XV to fully develop from ice $\beta-\mathrm{XV}$, but 2.5 minutes in the temperature interval 103128 K are not. These calorimetric findings suggest that the transition sequence is a three-step process, as follows: In the first step ice $\beta$-XV transforms to ice $\mathrm{VI}^{\ddagger}$ at 103 K , causing the first endotherm. In the second step, between about 103 K and 128 K ice $\mathrm{VI}^{\ddagger}$ develops to the more stable ordered ice XV , where kinetics are the limitation. This sequence via ice $\mathrm{VI}^{\ddagger}$ describes how the first H -order-to-H-order transition in ice physics takes place. In the third step ice XV disorders to produce metastable ice VI.

Now let's move on to the procedure for ruling out space groups using Bilbao. In order to check for the validity of a given space group, the O -atoms were reflected through each symmetry element. Within valid space groups these reflections have to match exactly with positions of other O atoms. This point is illustrated in Supplementary Figures 1 and 2. In Supplementary Figure 1, the Oatom circled in green was analysed for space group Pcc2. Its first reflected position can be obtained by reflection through the y -z-plane and a subsequent translation of half the lattice vector along z -direction (from -0.1 to 0.4 ). As we can find another O atom (circled in black) on that position, this symmetry operation is valid. Analogously, we can verify the reflection through the x-z-plane and the subsequent translation. Also, the twofold axis can be confirmed easily. In Supplementary Figure 2, we can find a different situation. Here, space group Pnc2 is considered and the O -atom circled in green has to be reflected through the y -z-plane and subsequently translated through half of the diagonal of the a-cface. At that position, however, we cannot find an oxygen atom. That can easily be visualized by the z positions as the greenly circled O -atom is, again, reflected from -0.1 to 0.4 , whereas the oxygen atom at the respective position within the $\mathrm{x}-\mathrm{y}$-plane is set at $\mathrm{z}=-0.4$.

Let's finally analyse the statistical quality of the fits for the best four models found in our Rietveld analysis:

|  | $w \mathrm{R}_{\mathrm{P}}(\%)$ | $\chi^{2}$ |
| :---: | :---: | :---: |
| $P \overline{4}$ | 2.38 | 3.278 |
| $P c c 2$ | 2.63 | 4.018 |
| $P 2_{1} 2_{1} 2$ | 2.82 | 4.588 |
| $P c a 2_{1}$ | 2.83 | 4.645 |

The weighted Rietveld powder R-factor, $w \mathrm{R}_{\mathrm{P}}$, is defined as:

$$
\begin{equation*}
\left(w R_{P}\right)^{2}=\sum_{i} w_{i}\left[y_{i}(o b s)-y_{i}(c a l c)\right]^{2} / \sum_{i} w_{i}\left[y_{i}(o b s)\right]^{2} \tag{1}
\end{equation*}
$$

where the sum is over $i$ observed points with observed and calculated intensities, $y_{i}(\mathrm{obs})$ and $y_{i}(\mathrm{calc})$, respectively, and a statistical weight $w_{i}$. The chi-squared $\left(\chi^{2}\right)$ is then defined as:

$$
\begin{equation*}
\chi^{2}=\left(w R_{P} / R_{\text {exp }}\right)^{2} \tag{2}
\end{equation*}
$$

The 'expected', or best-possible, or R-factor $\left(\mathrm{R}_{\mathrm{exp}}\right)$ is:

$$
\begin{equation*}
\left(R_{\text {exp }}\right)^{2}=(N-P+C) / \sum_{i} w_{i}\left[y_{i}(o b s)\right]^{2} \tag{3}
\end{equation*}
$$

Where $N$ is the total number of observations, $P$ is the number of refined parameters and $C$ is the number of applied constraints.

## Supplementary Figures

Supplementary Figure 1: Space group analysis for Pcc2. $\sqrt{ } 2 x \sqrt{ } 2 x 1$ super-cell of the ice VI unit cell containing the reflection conditions of Pcc2: full circles represent $O$-atoms, empty circles $H$-sites, different colours different $H$-sublattices and white numbers the $z$-positions of the O -atoms.


Supplementary Figure 2: Space group analysis for Pnc2. $\sqrt{ } 2 x \sqrt{ } 2 x 1$ super-cell of the ice VI unit cell containing the reflection conditions of Pnc2: full circles represent $O$-atoms, empty circles $H$-sites, different colours different $H$-sublattices and white numbers the $z$-positions of the $O$-atoms.


## Supplementary Figure 3: Occupancies of the H-sites according to Pcc2 symmetry operators.



Supplementary Figure 4: Rietveld refinement of Pcc2 model. Neutron powder diffraction data for ice XIX (red circular symbols) and the fit calculated from the refined Pcc2 model (green line), with the background subtracted. The difference between the model and data is represent by the purple line underneath the diffraction pattern and the positions of the Bragg peaks are indicated by vertical black tick marks. The inset shows the region between 1.7 and $2.45 \AA$-spacing where peaks appear that are permitted by the ice VI cell metric but have little or no intensity in ice VI (*) and where super-lattice peaks are observed ( $\dagger$ ).


Supplementary Figure 5: Unit-cell of ice XIX (refined Pcc2 model). View along the c-axis, showing the atom labelling scheme used in Supplementary Table 1.


Supplementary Figure 6: Rietveld refinement of $\boldsymbol{P 2}_{2_{1} \mathbf{2}_{1} \mathbf{2}}$ model. Neutron powder diffraction data for ice XIX (red circular symbols) and the fit calculated from the refined $P 2_{1} 2_{1} 2$ model (green line), with the background subtracted. The difference between the model and data is represent by the purple line underneath the diffraction pattern and the positions of the Bragg peaks are indicated by vertical black tick marks. The inset shows the region between 1.7 and $2.45 \AA$ d-spacing where peaks appear that are permitted by the ice VI cell metric but have little or no intensity in ice VI (*) and where superlattice peaks are observed ( $\uparrow$ ).


Supplementary Figure 7: Unit-cell of ice XIX (refined P2 $\mathbf{2 1}_{2} 2$ model). View along the $c$-axis, showing the atom labelling scheme used in Supplementary Table 2.


Supplementary Figure 8: Rietveld refinement of Pca2 $\mathbf{1 m}_{1}$ model. Neutron powder diffraction data for ice XIX (red circular symbols) and the fit calculated from the refined Pca2 $2_{1}$ model (green line), with the background subtracted. The difference between the model and data is represent by the purple line underneath the diffraction pattern and the positions of the Bragg peaks are indicated by vertical black tick marks. The inset shows the region between 1.7 and $2.45 \AA d$-spacing where peaks appear that are permitted by the ice VI cell metric but have little or no intensity in ice VI (*) and where super-lattice peaks are observed ( $\dagger$ ).


Supplementary Figure 9: Unit-cell of ice XIX (refined Pca2 $\mathbf{1}_{1}$ model). View along the $c$-axis, showing the atom labelling scheme used in Supplementary Table 3.


## Supplementary Tables

Supplementary Table 1: Refined structure using the Pcc2 model
$a=8.84253$ (7) $\AA$
$b=8.82654$ (7) $\AA$
$c=5.75559(5) \AA$
$\mathrm{V}=449.218(5) \AA^{3}$

| Atom | Wyckoff <br> position | $x$ | $y$ | $z$ | Occ. | $\mathrm{U}_{\text {iso }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |  |
| O1 | $4 e$ | 0.25 | 0.25 | 0.25 | 1.0 | 0.0081 |
| O2 | $4 e$ | $0.1107(2)$ | $0.1074(2)$ | $0.8770(3)$ | 1.0 | 0.0081 |
| O3 | $4 e$ | $0.1092(2)$ | $0.3881(2)$ | $0.6322(3)$ | 1.0 | 0.0081 |
| O4 | $4 e$ | $0.3892(2)$ | $0.3884(2)$ | $0.8695(3)$ | 1.0 | 0.0081 |
| O5 | $4 e$ | $0.3888(2)$ | $0.1073(2)$ | $0.6259(3)$ | 1.0 | 0.0081 |
| D11 | $4 e$ | $0.1869(4)$ | $0.1933(5)$ | $0.1555(6)$ | $0.545(3)$ | 0.0219 |
| D12 | $4 e$ | $0.1931(4)$ | $0.3143(4)$ | $0.3424(6)$ | $0.579(2)$ | 0.0219 |
| D13 | $4 e$ | $0.3036(5)$ | $0.3046(5)$ | $0.1400(6)$ | $0.471(3)$ | 0.0219 |
| D14 | $4 e$ | $0.3066(5)$ | $0.1981(6)$ | $0.3595(6)$ | $0.383(3)$ | 0.0219 |
| D21 | $4 e$ | $0.0356(4)$ | $0.0340(4)$ | $0.8910(14)$ | 0.5 | 0.0219 |
| D22 | $4 e$ | $0.0814(10)$ | $0.1914(4)$ | $0.7917(9)$ | $0.271(2)$ | 0.0219 |
| D23 | $4 e$ | $0.1497(7)$ | $0.1425(7)$ | $0.0161(5)$ | $0.455(3)$ | 0.0219 |
| D24 | $4 e$ | $0.1994(3)$ | $0.0774(6)$ | $0.8042(7)$ | $0.783(1)$ | 0.0219 |
| D31 | $4 e$ | $0.0359(4)$ | $0.4632(4)$ | $0.6165(15)$ | 0.5 | 0.0219 |
| D32 | $4 e$ | $0.1967(4)$ | $0.4293(7)$ | $0.6941(10)$ | $0.357(3)$ | 0.0219 |
| D33 | $4 e$ | $0.1389(8)$ | $0.3537(8)$ | $0.4875(5)$ | $0.422(2)$ | 0.0219 |
| D34 | $4 e$ | $0.0859(7)$ | $0.3036(3)$ | $0.7216(6)$ | $0.729(2)$ | 0.0219 |
| D41 | $4 e$ | $0.4660(4)$ | $0.4601(4)$ | $0.8609(13)$ | 0.5 | 0.0219 |
| D42 | $4 e$ | $0.4042(12)$ | $0.3005(4)$ | $0.7845(8)$ | $0.331(3)$ | 0.0219 |
| D43 | $4 e$ | $0.3639(8)$ | $0.3475(6)$ | $0.0121(5)$ | $0.529(3)$ | 0.0219 |
| D44 | $4 e$ | $0.3032(3)$ | $0.4039(8)$ | $0.7798(6)$ | $0.644(3)$ | 0.0219 |
| D51 | $4 e$ | $0.4647(4)$ | $0.0344(4)$ | $0.6326(15)$ | 0.5 | 0.0219 |
| D52 | $4 e$ | $0.3065(4)$ | $0.0847(13)$ | $0.7192(8)$ | $0.217(1)$ | 0.0219 |
| D53 | $4 e$ | $0.3496(6)$ | $0.1451(6)$ | $0.4887(5)$ | $0.617(3)$ | 0.0219 |
| D54 | $4 e$ | $0.4189(6)$ | $0.1939(3)$ | $0.7048(7)$ | $0.669(3)$ | 0.0219 |

Supplementary Table 2: Refined structure using the $P 2_{2_{l}} 2$ model.

$$
\begin{aligned}
& a=8.84210(11) \AA \\
& b=8.82682(12) \AA \\
& c=5.75549(6) \AA \\
& \mathrm{V}=449.202(4) \AA \AA^{3}
\end{aligned}
$$

| Atom | Wyckoff <br> position | $x$ | $y$ | $z$ | Occ. | $\mathrm{U}_{\text {iso }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |  |
| O1 | $2 a$ | 0.00 | 0.00 | $0.503(3)$ | 1.0 | 0.0098 |
| O2 | $2 b$ | 0.00 | 0.50 | $0.007(2)$ | 1.0 | 0.0098 |
| O3 | $4 c$ | $0.139(1)$ | $0.146(1)$ | $0.132(1)$ | 1.0 | 0.0098 |
| O4 | $4 c$ | $0.138(1)$ | $0.355(1)$ | $0.612(1)$ | 1.0 | 0.0098 |
| O5 | $4 c$ | $0.361(1)$ | $0.364(1)$ | $0.123(1)$ | 1.0 | 0.0098 |
| O6 | $4 c$ | $0.363(1)$ | $0.142(1)$ | $0.622(1)$ | 1.0 | 0.0098 |
| D11 | $4 c$ | $0.061(3)$ | $0.052(3)$ | $0.399(3)$ | $0.384(1)$ | 0.0230 |
| D12 | $4 c$ | $0.436(2)$ | $0.444(2)$ | $0.404(2)$ | $0.616(1)$ | 0.0230 |
| D21 | $4 c$ | $0.056(2)$ | $0.437(2)$ | $0.912(3)$ | $0.400(1)$ | 0.0230 |
| D22 | $4 c$ | $-0.052(1)$ | $0.440(1)$ | $0.111(2)$ | $0.600(1)$ | 0.0230 |
| D31 | $4 c$ | $0.216(1)$ | $0.216(1)$ | $0.143(2)$ | $0.65(1)$ | 0.0230 |
| D32 | $4 c$ | $0.165(4)$ | $0.062(2)$ | $0.044(5)$ | $0.28(1)$ | 0.0230 |
| D33 | $4 c$ | $0.101(2)$ | $0.104(2)$ | $0.267(2)$ | $0.62(1)$ | 0.0230 |
| D34 | $4 c$ | $0.060(1)$ | $0.190(2)$ | $0.050(3)$ | $0.46(2)$ | 0.0230 |
| D41 | $4 c$ | $0.214(2)$ | $0.283(2)$ | $0.613(3)$ | $0.47(2)$ | 0.0230 |
| D42 | $4 c$ | $0.055(1)$ | $0.346(2)$ | $0.512(2)$ | $0.51(1)$ | 0.0230 |
| D43 | $4 c$ | $0.103(2)$ | $0.390(2)$ | $0.754(2)$ | $0.600(1)$ | 0.0230 |
| D44 | $4 c$ | $0.159(3)$ | $0.438(2)$ | $0.519(3)$ | $0.43(1)$ | 0.0230 |
| D51 | $4 c$ | $0.280(1)$ | $0.297(2)$ | $0.132(4)$ | $0.35(1)$ | 0.0230 |
| D52 | $4 c$ | $0.333(2)$ | $0.452(1)$ | $0.048(2)$ | $0.72(1)$ | 0.0230 |
| D53 | $4 c$ | $0.399(3)$ | $0.390(3)$ | $0.268(2)$ | $0.384(1)$ | 0.0230 |
| D54 | $4 c$ | $0.445(1)$ | $0.323(2)$ | $0.049(3)$ | $0.55(2)$ | 0.0230 |
| D61 | $4 c$ | $0.293(1)$ | $0.220(1)$ | $0.614(3)$ | $0.53(2)$ | 0.0230 |
| D62 | $4 c$ | $0.457(1)$ | $0.155(2)$ | $0.554(2)$ | $0.49(1)$ | 0.0230 |
| D63 | $4 c$ | $0.403(3)$ | $0.111(3)$ | $0.762(2)$ | $0.400(1)$ | 0.0230 |
| D64 | $4 c$ | $0.338(2)$ | $0.055(1)$ | $0.540(3)$ | $0.58(1)$ | 0.0230 |

Supplementary Table 3: Refined structure using the $P \subset a 2_{1}$ model
$a=5.75560(6) \AA$
$b=8.84224(10) \AA$
$c=8.82667(11) \AA$
$\mathrm{V}=449.211(4) \AA$

| Atom | Wyckoff <br> position | $x$ | $y$ | $z$ | Occ. | $\mathrm{U}_{\text {iso }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |  |
| O1 | $4 a$ | $0.505(2)$ | $0.253(1)$ | $0.003(1)$ | 1.0 | 0.0095 |
| O2 | $4 a$ | $0.620(2)$ | $0.115(1)$ | $0.358(1)$ | 1.0 | 0.0095 |
| O3 | $4 a$ | $0.116(2)$ | $0.113(1)$ | $0.142(1)$ | 1.0 | 0.0095 |
| O4 | $4 a$ | $0.875(2)$ | $0.394(1)$ | $0.140(1)$ | 1.0 | 0.0095 |
| O5 | $4 a$ | $0.371(2)$ | $0.389(1)$ | $0.357(1)$ | 1.0 | 0.0095 |
| D11 | $4 a$ | $0.896(2)$ | $0.188(1)$ | $0.451(2)$ | $0.65(1)$ | 0.0223 |
| D12 | $4 a$ | $0.422(3)$ | $0.190(3)$ | $0.067(2)$ | $0.34(1)$ | 0.0223 |
| D13 | $4 a$ | $0.602(2)$ | $0.307(2)$ | $0.066(2)$ | $0.58(2)$ | 0.0223 |
| D14 | $4 a$ | $0.110(3)$ | $0.304(2)$ | $0.449(2)$ | $0.43(2)$ | 0.0223 |
| D21 | $4 a$ | $0.623(3)$ | $0.044(2)$ | $0.281(1)$ | $0.58(2)$ | 0.0223 |
| D22 | $4 a$ | $0.530(3)$ | $0.103(2)$ | $0.445(2)$ | $0.52(2)$ | 0.0223 |
| D23 | $4 a$ | $0.768(2)$ | $0.146(4)$ | $0.387(3)$ | $0.35(1)$ | 0.0223 |
| D24 | $4 a$ | $0.541(3)$ | $0.205(1)$ | $0.343(2)$ | $0.56(2)$ | 0.0223 |
| D31 | $4 a$ | $0.111(3)$ | $0.033(2)$ | $0.210(3)$ | $0.42(2)$ | 0.0223 |
| D32 | $4 a$ | $0.034(3)$ | $0.200(2)$ | $0.166(3)$ | $0.44(1)$ | 0.0223 |
| D33 | $4 a$ | $0.259(2)$ | $0.149(2)$ | $0.110(2)$ | $0.66(1)$ | 0.0223 |
| D34 | $4 a$ | $0.031(4)$ | $0.091(2)$ | $0.055(2)$ | $0.49(2)$ | 0.0223 |
| D41 | $4 a$ | $0.866(3)$ | $0.474(2)$ | $0.208(2)$ | $0.48(2)$ | 0.0223 |
| D42 | $4 a$ | $0.941(3)$ | $0.428(2)$ | $0.051(1)$ | $0.54(2)$ | 0.0223 |
| D43 | $4 a$ | $0.735(3)$ | $0.355(3)$ | $0.106(3)$ | $0.42(2)$ | 0.0223 |
| D44 | $4 a$ | $0.940(3)$ | $0.306(1)$ | $0.179(2)$ | $0.57(1)$ | 0.0223 |
| D51 | $4 a$ | $0.368(3)$ | $0.463(2)$ | $0.282(2)$ | $0.52(2)$ | 0.0223 |
| D52 | $4 a$ | $0.462(4)$ | $0.306(2)$ | $0.329(2)$ | $0.45(2)$ | 0.0223 |
| D53 | $4 a$ | $0.236(2)$ | $0.350(2)$ | $0.397(2)$ | $0.57(2)$ | 0.0223 |
| D54 | $4 a$ | $0.475(3)$ | $0.420(2)$ | $0.431(2)$ | $0.46(2)$ | 0.0223 |

